

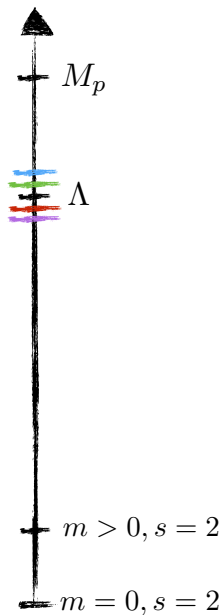
# CONSTRAINTS ON SPIN-2 INTERACTIONS

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[1712.10020](#), [1804.08686](#), [1806.xxxxx](#)

Amplitudes School, June 12th, 2018

# MOTIVATION

- ▶ Can there exist a large- $N$  QCD theory with an isolated massive spin-2 glueball as the lightest state?
- ▶ Is there a bound on the gap to the next lightest state?



# CAUSALITY CONSTRAINTS

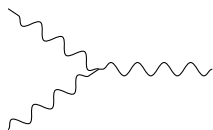
- ▶ To address the first question, we calculate the Shapiro time delay/advance in this theory [CEMZ, 2014](#).
- ▶ This is captured by the Eikonal scattering amplitude:

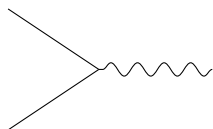
$$i\mathcal{M}_{\text{eik}}(s, t) = 2s \int d^2\vec{b} e^{i\vec{q}\cdot\vec{b}} \left( e^{i\delta(s, \vec{b})} - 1 \right).$$

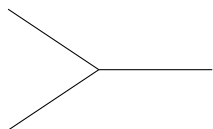
- ▶ The Eikonal phase  $\delta$  depends only on on-shell cubic vertices.
- ▶ A time advance,  $\delta < 0$ , would imply that new physics is needed around the mass scale  $m$  to restore causality.

# CONSTRAINING CUBICS

- ▶ Prohibiting time advances gives the following cubic vertices:


$$= \mathcal{V}_{\text{EH}},$$

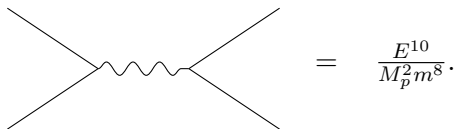

$$= \mathcal{V}_{\text{EH}},$$


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$$\mathcal{V}_{\text{EH}} = \frac{2i}{M_p} [(\epsilon_1 \cdot \epsilon_2)(\epsilon_3 \cdot p_1) + (\epsilon_1 \cdot \epsilon_3)(\epsilon_2 \cdot p_3) + (\epsilon_2 \cdot \epsilon_3)(\epsilon_1 \cdot p_2)]^2.$$

## CONSTRAINING THE GAP

- ▶ Tree amplitudes violate unitarity at  $(m^4 M_p)^{1/5}$ , so the EFT cutoff is below this scale.


$$= \frac{E^{10}}{M_p^2 m^8}.$$

- ▶ Contact terms can cancel the bad high-energy behaviour and raise the cutoff, increasing the gap.
- ▶ To find the maximum cutoff, calculate the general four-point amplitude consistent with locality, unitarity, Lorentz invariance, gauge invariance, and crossing symmetry.
- ▶ This gives  $(m^2 M_p)^{1/3}$  as the highest cutoff.

# CONCLUSION

- ▶ Can there exist a large- $N$  theory with an isolated massive spin-2 glueball as the lightest state?  
Perhaps! Causality implies that on-shell cubic vertices must be of EH form.
- ▶ Is there a bound on the gap to the next lightest state?  
Perturbative unitarity implies that new states must enter by  $\Lambda = (m^2 M_p)^{1/3}$ .
- ▶ Can generalize beyond this simple example to other particle spectra, e.g. higher spins, Higgs-like particles.

