

# 4-GLUON AMPLITUDE AT 3-LOOPS IN QCD

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# PROLOGUE

- Scattering amplitude is the basic building block in QFT
- More precise theoretical predictions are crying need at the LHC
- Necessary to reveal the underlying structures in QFT
- We focus on a **4-gluon process** involving on-shell partons in QCD

$$g(p_1) + g(p_2) + g(p_3) + g(p_4) \rightarrow 0$$

- Ingredient for the di-jet production
- State-of-the-art
  - 1-loop: **Ellis, Sexton '86**
  - 2-loop Helicity amplitude: **Bern, Dixon, Kosower '00**
  - 2-loop full amplitude: **Glover, Oleari, Yeomans '01**
- **Next challenging goal: go for 3-loop!**
- First attempt in N=4 by **Henn, Mistlberger in '16**
- We address this **4-point amplitude at 3-loop for the first time in QCD**

Our goal



# PROCESS OF INTEREST

- We consider

$$g(p_1) + g(p_2) + g(p_3) + g(p_4) \rightarrow 0 \quad \text{on-shell in QCD}$$

with arbitrary number of light fermions in fund represent of  $SU(N)$  gauge group

- **Goal:** Compute the amplitude and helicity amplitudes in **planar limit**

Include all the terms satisfying  $n_f^{a_1} N^{a_2} a_s^{a_3}$  with  $a_1 + a_2 = a_3$

- One approach: decompose the amplitude into linearly independent tensor structures
- **138** Tensorial Structures

$$\text{constraints} \left\{ \begin{array}{l} \text{Transversality Condition} \quad \epsilon_i \cdot p_i = 0 \\ \text{Gauge Fixed Tensor} \quad \epsilon_1 \cdot p_2 = \epsilon_2 \cdot p_3 = \epsilon_3 \cdot p_4 = \epsilon_4 \cdot p_1 = 0 \end{array} \right.$$

Number of tensorial structures reduce to **10**

- All plus amplitude

$$|\mathcal{M}_{++++}\rangle = \sum_{i=1}^{10} A_i T_{i,++++} = \frac{1}{4\langle 12\rangle\langle 23\rangle\langle 34\rangle\langle 41\rangle} \left\{ -4(A_1 + A_3)st - 4A_2u^2 - 2A_4stu + 2A_5su^2 - 2A_6stu \right. \\ \left. -2A_7stu + 2A_8tu^2 - 2A_9stu - A_{10}stu^2 \right\}$$

# COMPUTATION

- Age-old Feynman diagrammatic approach, however, lots of challenges!
- Huge expressions, complicated reductions!
- Feynman diagrams using Qgraf: 48723 three-loop, 39k planar topologies
- Discard non-planar diagrams by removing sub-leading colors.
- Cross-checked using REDUZE2 using the liberty of shifting loop momenta
- In-house FORM code: convert Qgraf raw output to FORM

Noguira '06

Monteuffel, Studerus '12

SU(N) color simplification

Lorentz algebra

Dirac algebra



in d-spacetime dimensions

Vermaseren

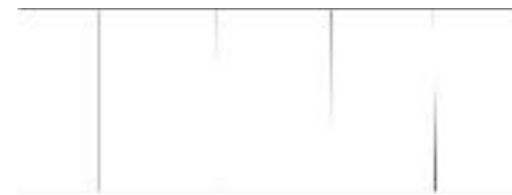
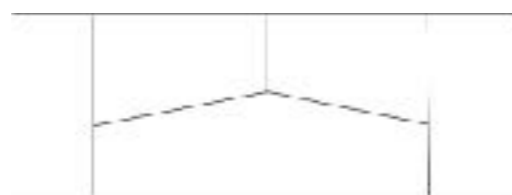
- Millions of 3-loop integrals with very high power of numerators: highest 6!

$$J[a_1, a_2, \dots, a_{15}] = \int \prod_{L=1,2,3} \left[ \frac{d^d k_L}{(2\pi)^d} \right] \frac{1}{D_1^{a_1} D_2^{a_2} \dots D_{15}^{a_{15}}}$$

D are the inverse of propagators

Most complicated

$$\sum |a_i| = 16$$



# COMPUTATION

- IBP Reduction

FIRE5.2 C++ along with LiteRed1.82

2-step reductions:

1. LiteRed along with Mint: symbolic rules & 89 MIs
2. Non-minimal set : Huge reduction file!
3. Reduce again using FIRE c++: 81 MIs
4. Minimal set: reduction file size reduced 1/10

Boels, Jin, Luo '18

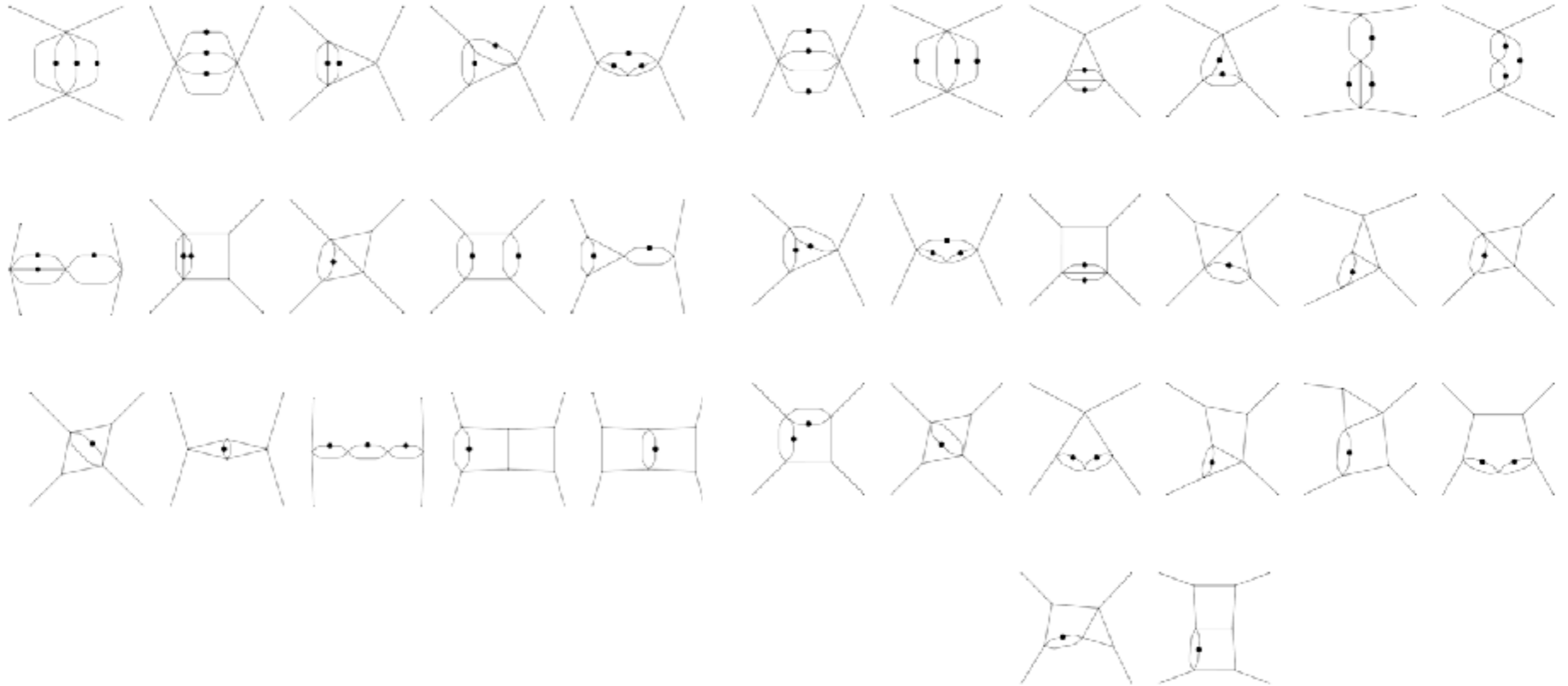
Advantage: substantial improvement of total reduction time

Comments: Non-minimal set of MIs is inefficient

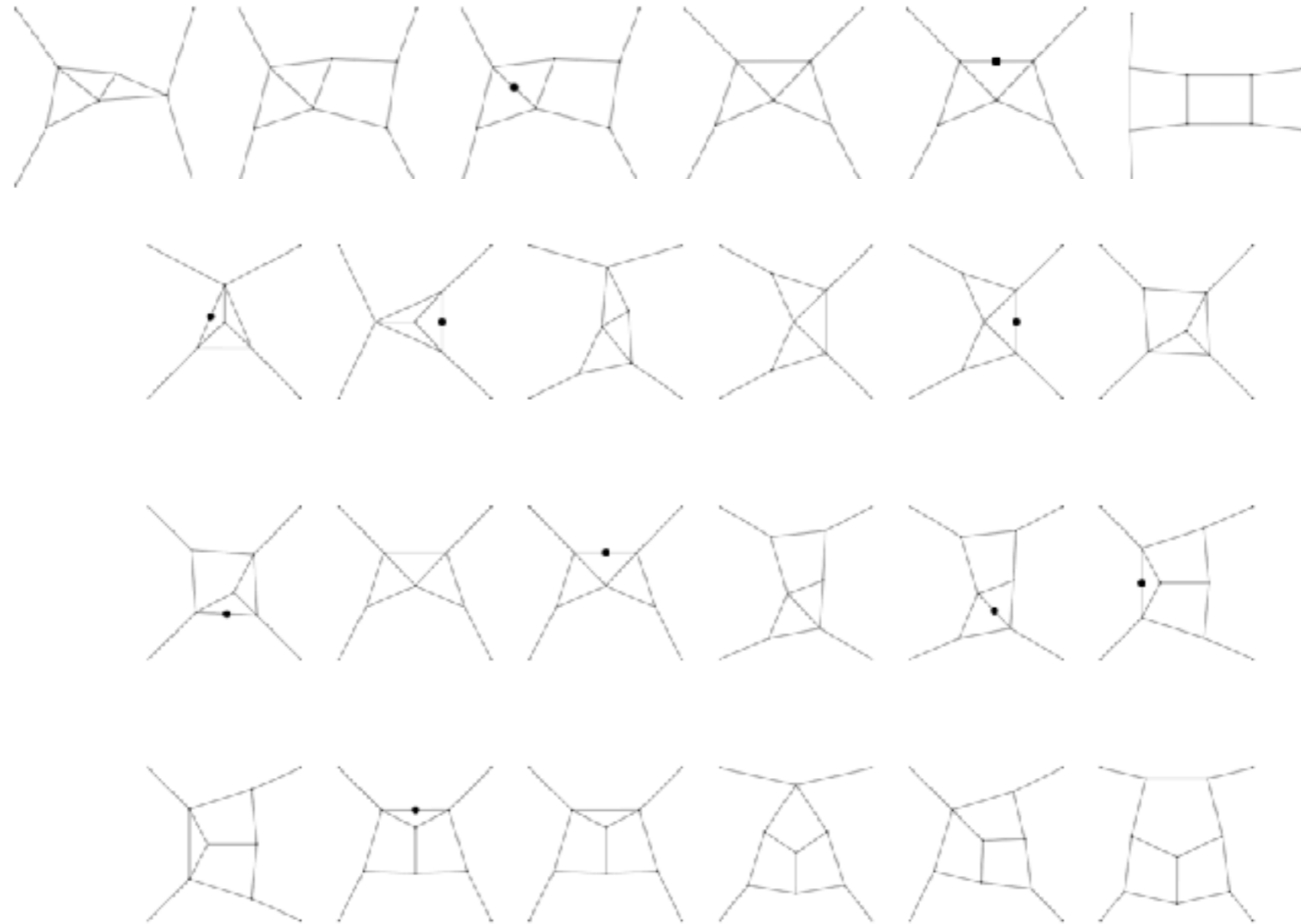
- 21 MIs with one double prop
- 2 MIs with one numerator

# MIs IN UT-BASIS

Henn, Smirnov, Smirnov '13



With Bubble Sub-diagrams



MIs without Bubble sub-integrals

- Remaining integrals are obtained by interchanging  $s$  &  $t$

# UV & IR

- Dimensional regularisation:  $d = 4 - 2\epsilon$
- UV structures of amplitude and all plus amplitude are different

$$\mathcal{M}^{(0)}(s, t, \epsilon) \neq 0 \quad \text{UV divergent 1-loop}$$

$$\mathcal{M}_{++++}^{(0)}(s, t, \epsilon) = 0 \quad \text{All plus amplitude vanishes at tree level} \quad \rightarrow \quad \text{UV Finite 1-loop}$$

UV renormalisation in  $\overline{\text{MS}}$ -bar

- The IR poles are universal and were first predicted by Catani up to 2-loop (except single pole)
- From SCET

$$\mathcal{M}^{fin}(s, t) = \lim_{\epsilon \rightarrow 0} \mathbf{Z}^{-1}(\epsilon, s, t) \mathcal{M}(\epsilon, s, t)$$

 UV finite

All the IR divergences are governed by the matrix  $\mathbf{Z}$

The all order solution:

$$\mathbf{Z}(\epsilon, s, t, \mu) = \mathcal{P} \exp \left[ \int_{\mu}^{\infty} \frac{d\mu'}{\mu'} \mathbf{\Gamma}(s, t, \mu') \right]$$

Becher & Neubert '09

Gardi & Magnea '09

Catani '98

Sterman, Yeomans '03



# RESULTS & CONCLUSIONS

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- Our 1 and 2-loop results agree with the universal IR structure and the literatures.
- We have checked the amplitude as well as + + + + amplitude
- They exhibit the necessary symmetry between s & t
- We have extended the 2-loop results to higher order in  $\epsilon$
- 3-loop reduction is done!
- Result in terms of Master integrals is at hand.
- Just waiting for the final result! Will be available soon!
- First step towards the full computation
- First ever attempt in QCD!

**Thank you!**

# PRELIMINARY RESULT (PARTIAL OF COURSE)

Coefficient of  $n_f^3 \text{Tr}(T^{a_1} T^{a_3} T^{a_4} T^{a_2})$  of the Amplitude (Unrenormalised)

$$\begin{aligned} & \frac{1}{\epsilon} \frac{1}{(1+x)^3} \left[ T_1 \left( \frac{8x^4}{27} + \frac{28x^3}{27} + \frac{4x^2}{3} + \frac{20x}{27} + \frac{4}{27} \right) + T_2 \left( -\frac{4x^3}{27} - \frac{4x^2}{27} + \frac{4x}{27} + \frac{4}{27} \right) + T_5 \left( \frac{8}{27t} - \frac{8x^2}{27t} \right) \right. \\ & \quad \left. + T_8 \left( \frac{8}{27t} - \frac{8x^2}{27t} \right) + T_{10} \left( \frac{16}{27t^2} - \frac{16x}{27t^2} \right) \right] \\ & + \frac{1}{(1+x)^3} \left[ T_1 \left\{ \left( -\frac{8x^4}{9} - \frac{28x^3}{9} - 4x^2 - \frac{20x}{9} - \frac{4}{9} \right) \left( H(\{0\}, x) + \log(-s) - 3 \right) \right\} \right. \\ & \quad + T_2 \left\{ \left( \frac{4x^3}{9} + \frac{4x^2}{9} - \frac{4x}{9} - \frac{4}{9} \right) \left( \log(-s) - 3 \right) \right\} \\ & \quad + T_5 \left\{ \left( \frac{8x^2}{9t} - \frac{8}{9t} \right) \left( \log(-s) - 3 \right) \right\} + T_8 \left\{ \left( \frac{8x^2}{9t} - \frac{8}{9t} \right) \left( \log(-s) - 3 \right) \right\} \\ & \quad \left. + T_{10} \left\{ \left( \frac{16x}{9t^2} - \frac{16}{9t^2} \right) \left( \log(-s) - 3 \right) \right\} \right] \end{aligned}$$

Up to an overall normalisation factor

$$x = \frac{u}{t}, t = 2p_1 \cdot p_3, u = 2p_2 \cdot p_3$$

Absolutely no check has been performed! Will be done soon!

# LARGE N LIMIT

- Amplitude is a tensor in Color space
- Can be decomposed in terms of traces of fundamental color generators of SU(N)
- Pure gauge amplitude can be expressed in terms of six color structures

$$C_1 = tr(1234) + tr(1432)$$

$$C_2 = tr(1243) + tr(1342)$$

$$C_3 = tr(1423) + tr(1324)$$

$$C_4 = tr(12)tr(34)$$

$$C_5 = tr(13)tr(24)$$

$$C_6 = tr(14)tr(23)$$

$$tr(T^{a_1}T^{a_2}T^{a_3}T^{a_4}) = tr(1234)$$

It can be written as

$$A_i^{(L)} = \sum_{\lambda=1}^3 \left\{ \sum_{k=0}^{\lfloor \frac{L}{2} \rfloor} N^{L-2k} A_{i,\lambda}^{(L),2k} \right\} C_\lambda + \sum_{\lambda=4}^6 \left\{ \sum_{k=0}^{\lfloor \frac{L-1}{2} \rfloor} N^{L-2k-1} A_{i,\lambda}^{(L),2k+1} \right\} C_\lambda$$

$A_{i,\lambda}^{(L),0}$  are leading order in N. Others are sub-leading

Only  $A_{i,\lambda}^{(L),0}$  contributes in the large N-limit

Only planar diagrams contribute: Planar limit



Our goal